

PARAMETRIC STUDY OF MODULATED LIGHT TRANSMISSION THROUGH NANOSCALE HOLE ARRAYS

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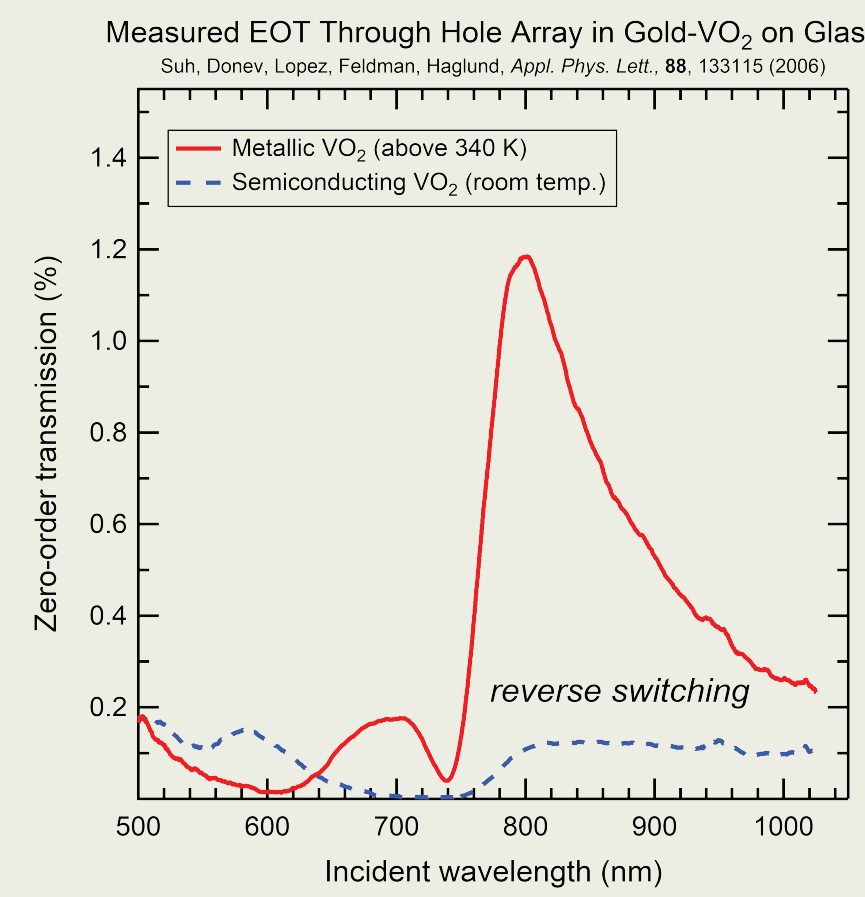
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THE EOT EFFECT

Before the discovery of *extraordinary optical transmission* (EOT) in 1998, which spurred the field of plasmonics, subwavelength holes in optically thick (~200 nm) metal films were expected to let through very little light. However, periodic nanohole arrays in plasmonic films, esp. silver (Ag) and gold (Au), can far exceed the expected weak transmission at visible and infrared (IR) wavelengths. The EOT effect is caused primarily by the excitation, interference, tunneling and scattering of *surface-plasmon polaritons* and *quasi-cylindrical waves*. Typical EOT spectra show asymmetric peaks preceded by sharp minima—a hallmark of *Fano-type interference* between resonant and continuum transmission channels.

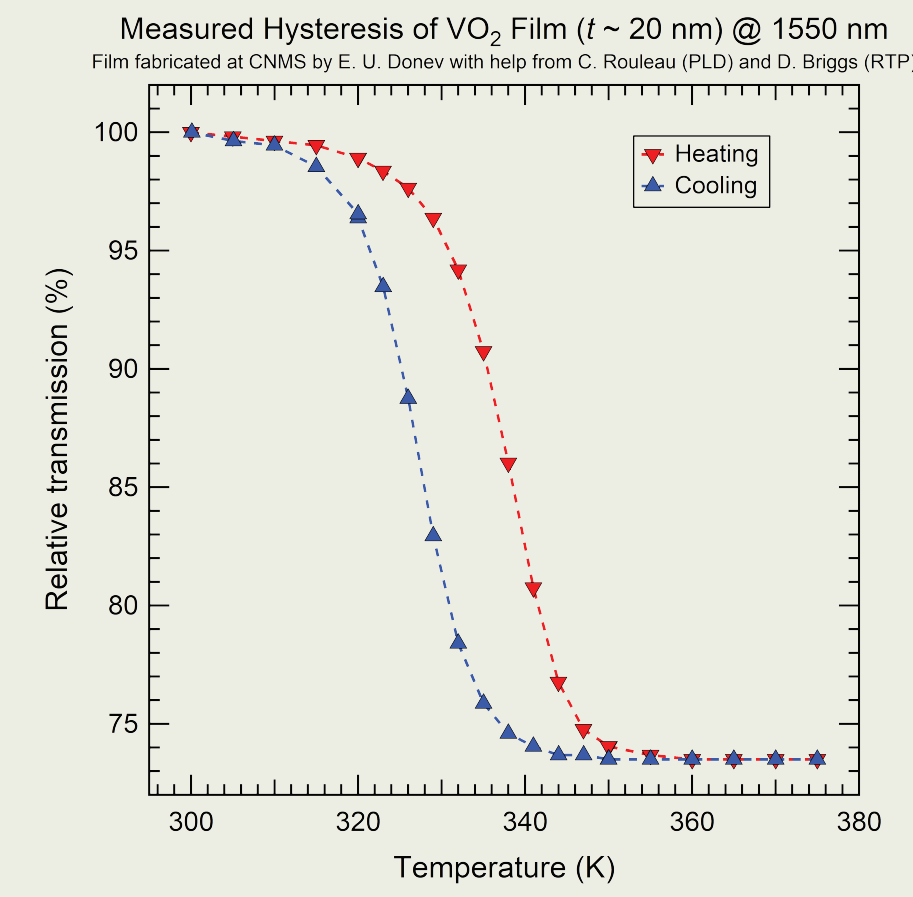
VO₂ AS EOT MODULATOR

The VO₂ phase transition is one of the few means to control the EOT effect dynamically. Previous experiments revealed a surprising effect, dubbed *reverse switching*, whereby Au+VO₂ and Ag+VO₂ nanohole arrays transmit more near-IR light when the VO₂ layer is in the metallic phase—opposite to what a plain VO₂ film does. The current study seeks to explore how the geometrical parameters of the array affect the reverse switching, and to understand its electromagnetic origins.



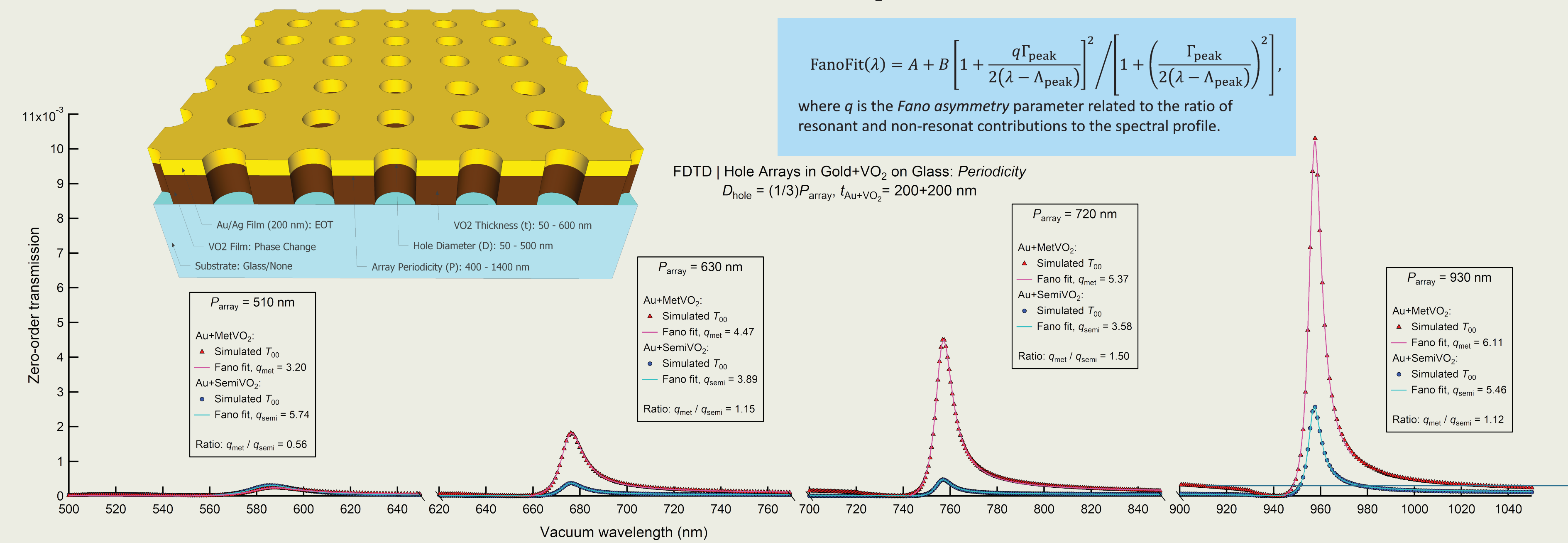
VO₂ PHASE TRANSITION

Vanadium dioxide (VO₂) is a correlated-electron material that undergoes phase changes induced by temperature or light. Above 340 K (67 °C), VO₂ switches from a semiconductor to a (“bad”) metal before switching back upon cooling with a hysteresis. Even a very thin plain (i.e., no holes) film is much opaquer to IR light in the metallic phase than in the semiconducting phase.



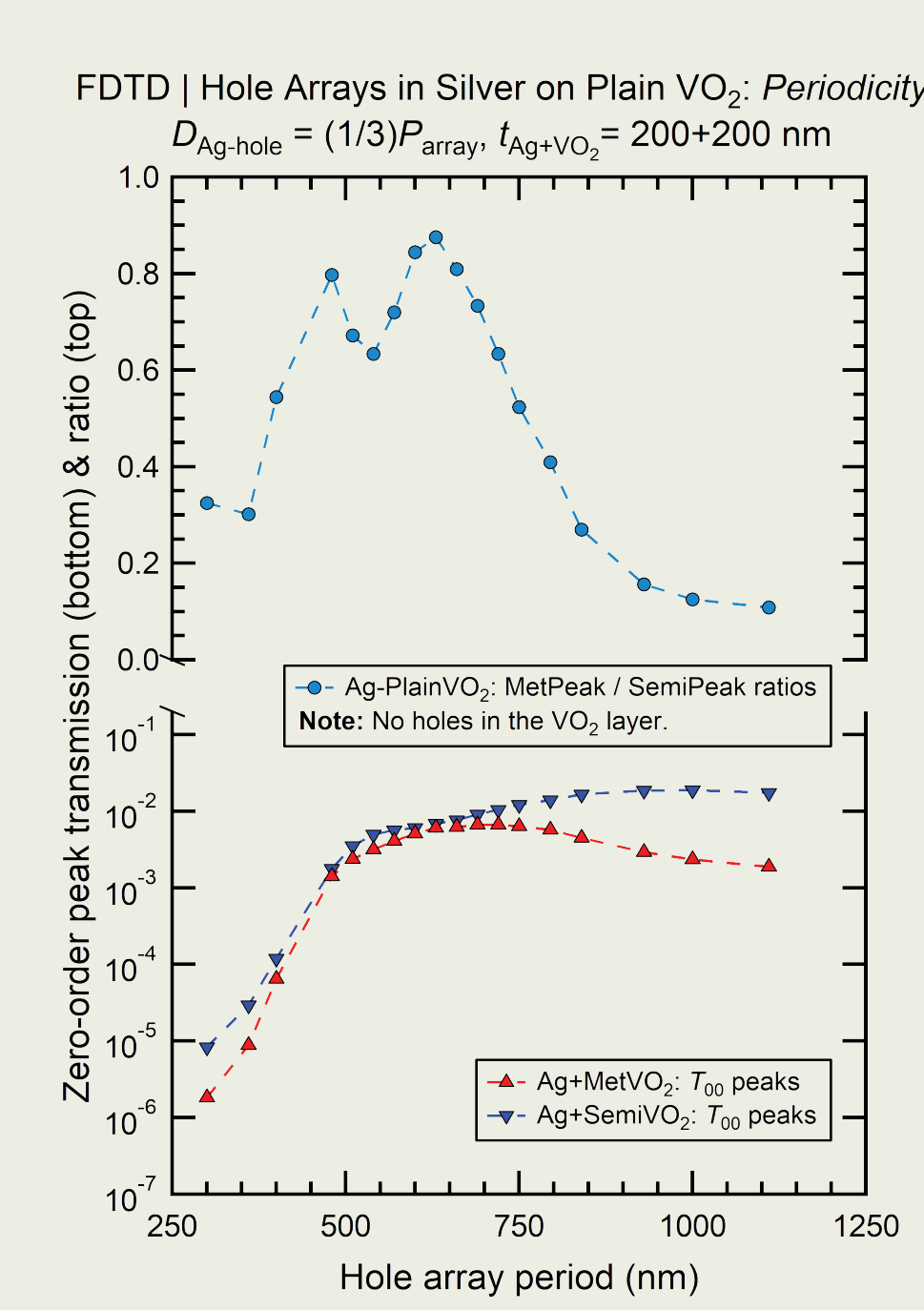
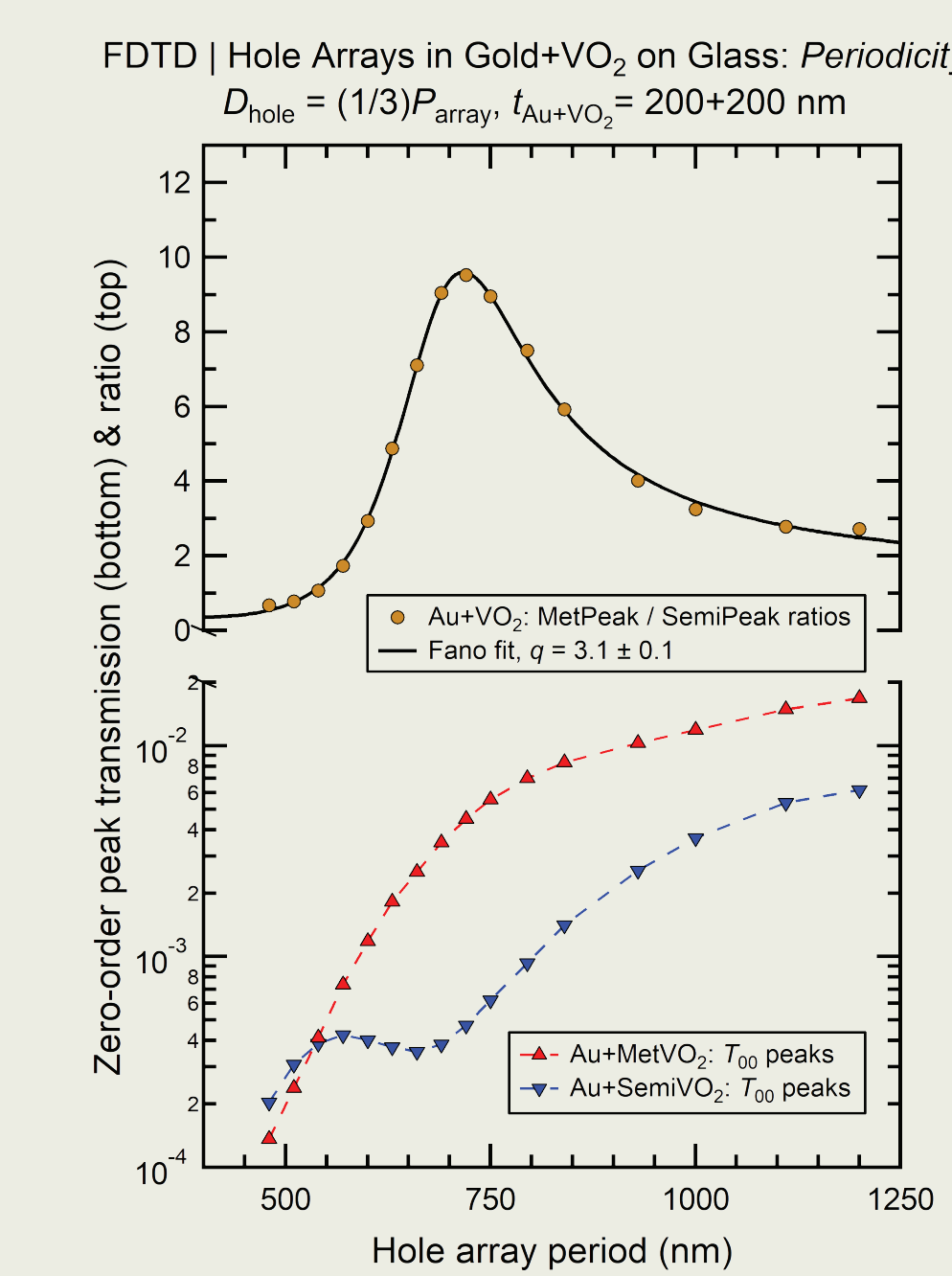
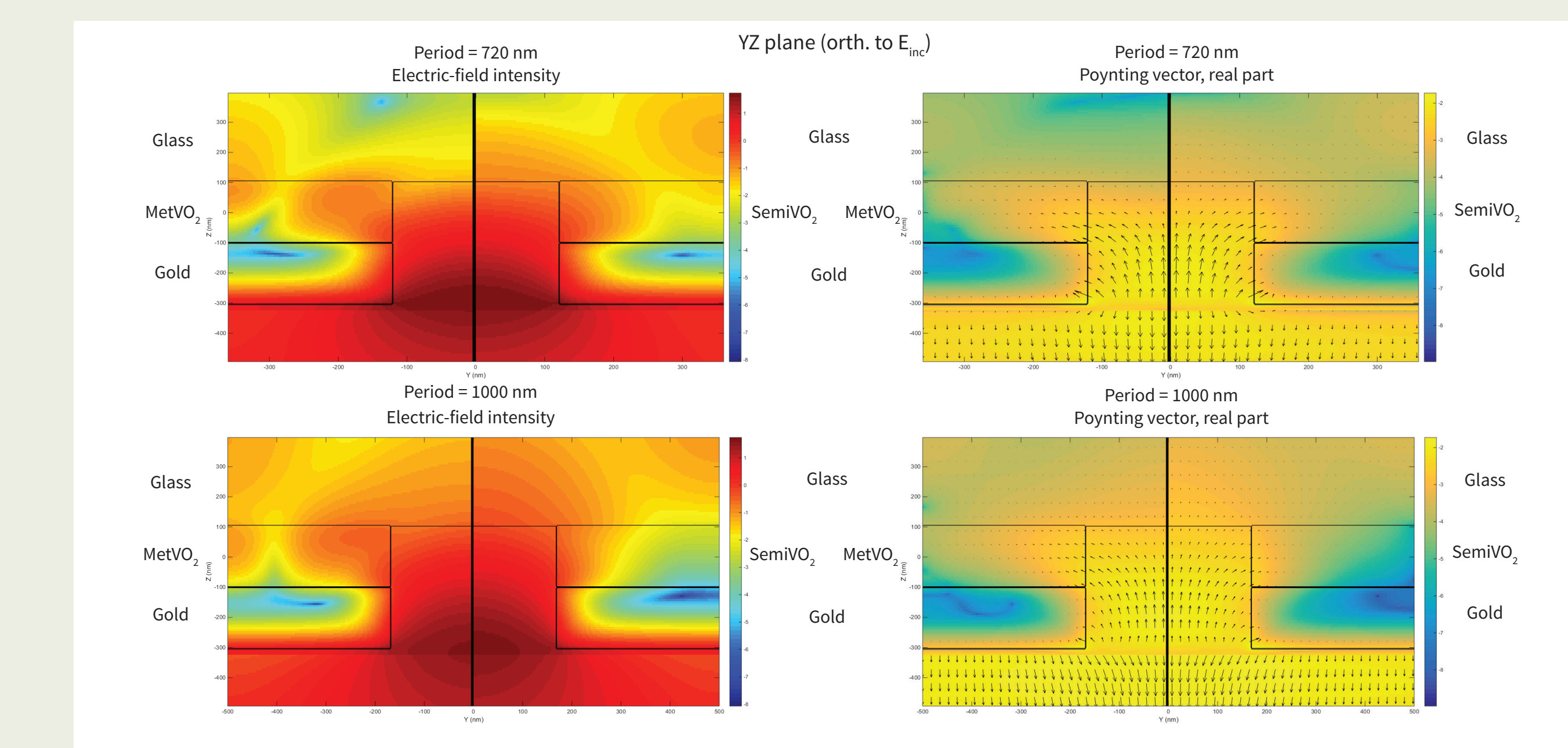
FULL-WAVE 3D NUMERICAL SIMULATIONS

Using the *finite-difference time-domain* (FDTD; Lumerical) method and *finite-element method* (FEM; COMSOL Multiphysics), we have simulated the optical response of nanohole arrays in bilayer films of Au+VO₂ and Ag+VO₂ while varying the periodicity (*P*), hole diameter (*D*), VO₂ thickness (*t*). In the process, we uncovered intriguing “meta-Fano” trends in the ratios of the transmission peaks for the two VO₂ phases, implying interactions between resonant and non-resonant contributions to the modulation mechanism. We suspect that a competition between field penetration and absorption causes a pronounced dip in the transmission of the semiconducting VO₂ within a narrow range of each geometrical parameter. We also identify a lossy Fabry-Perot resonance with variation in VO₂ thickness.



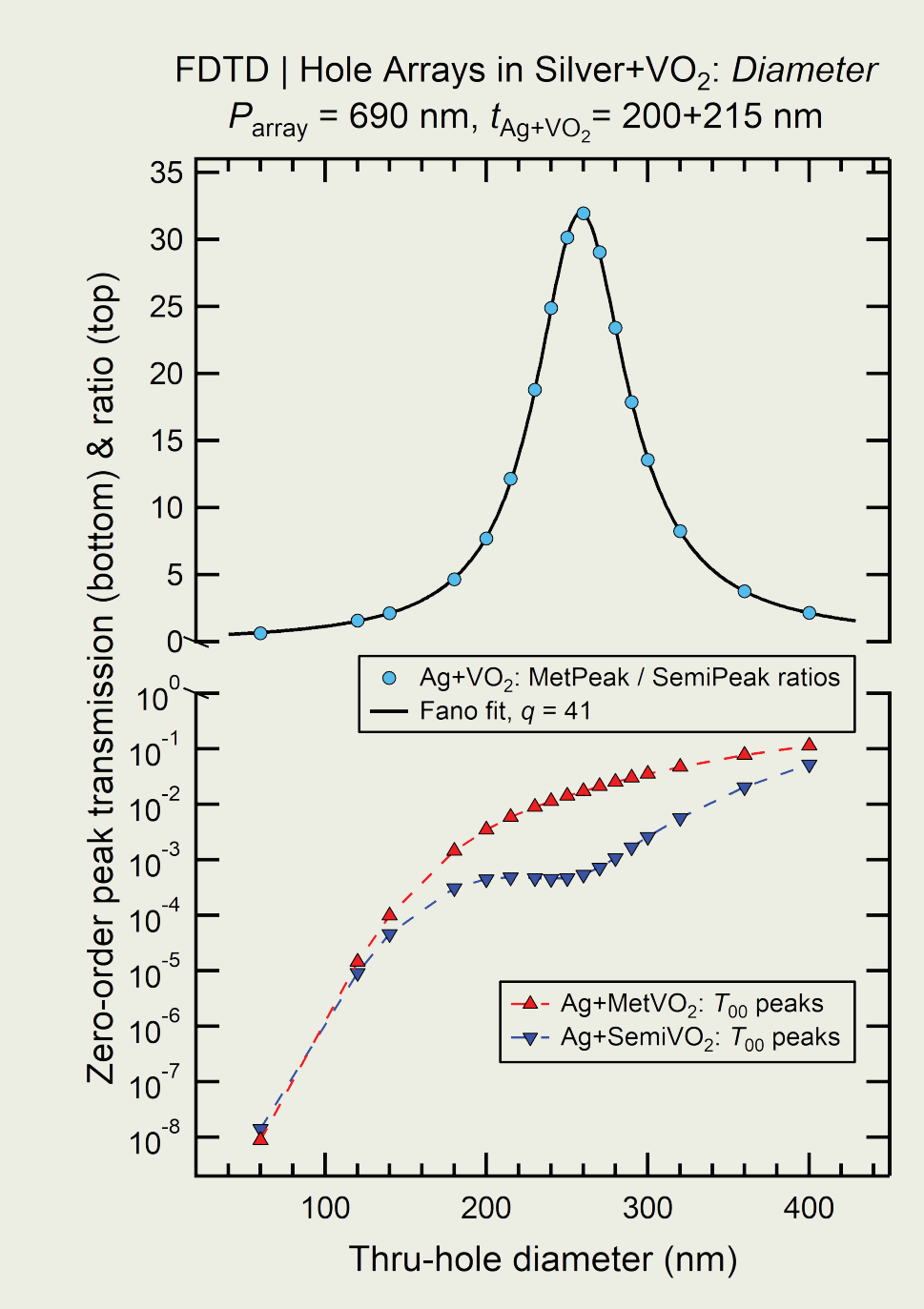
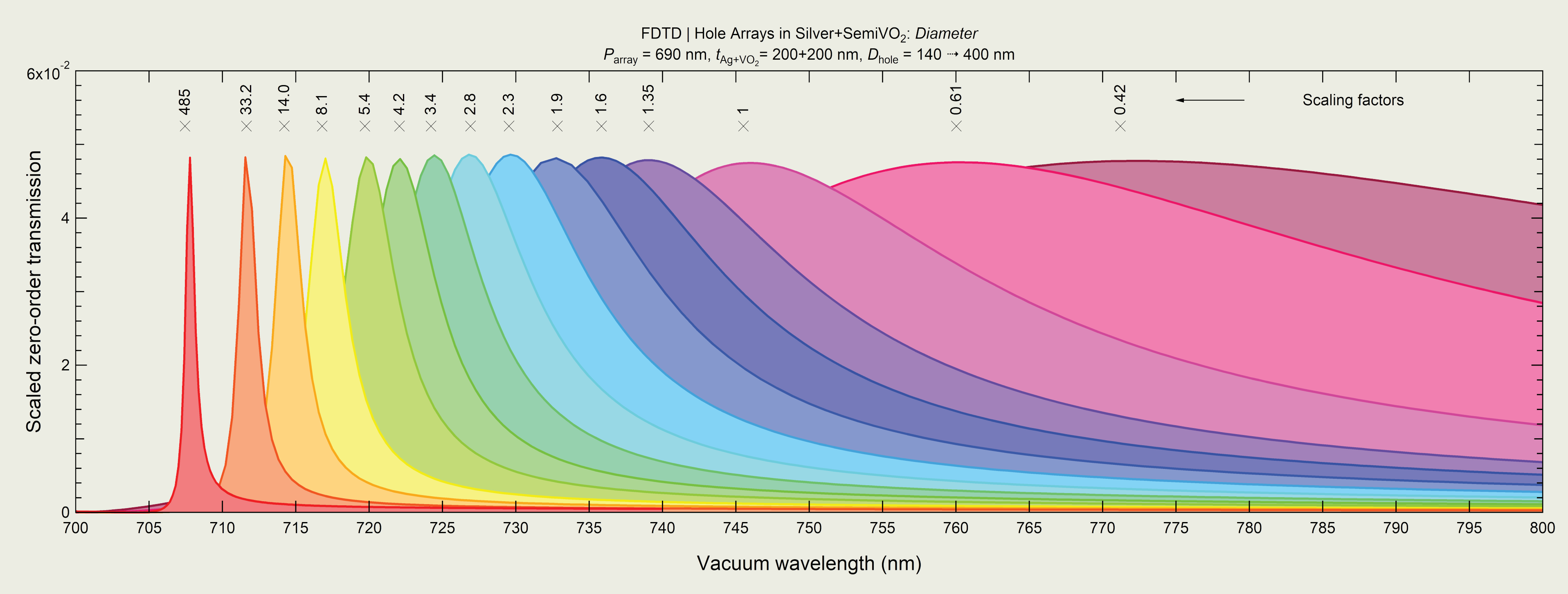
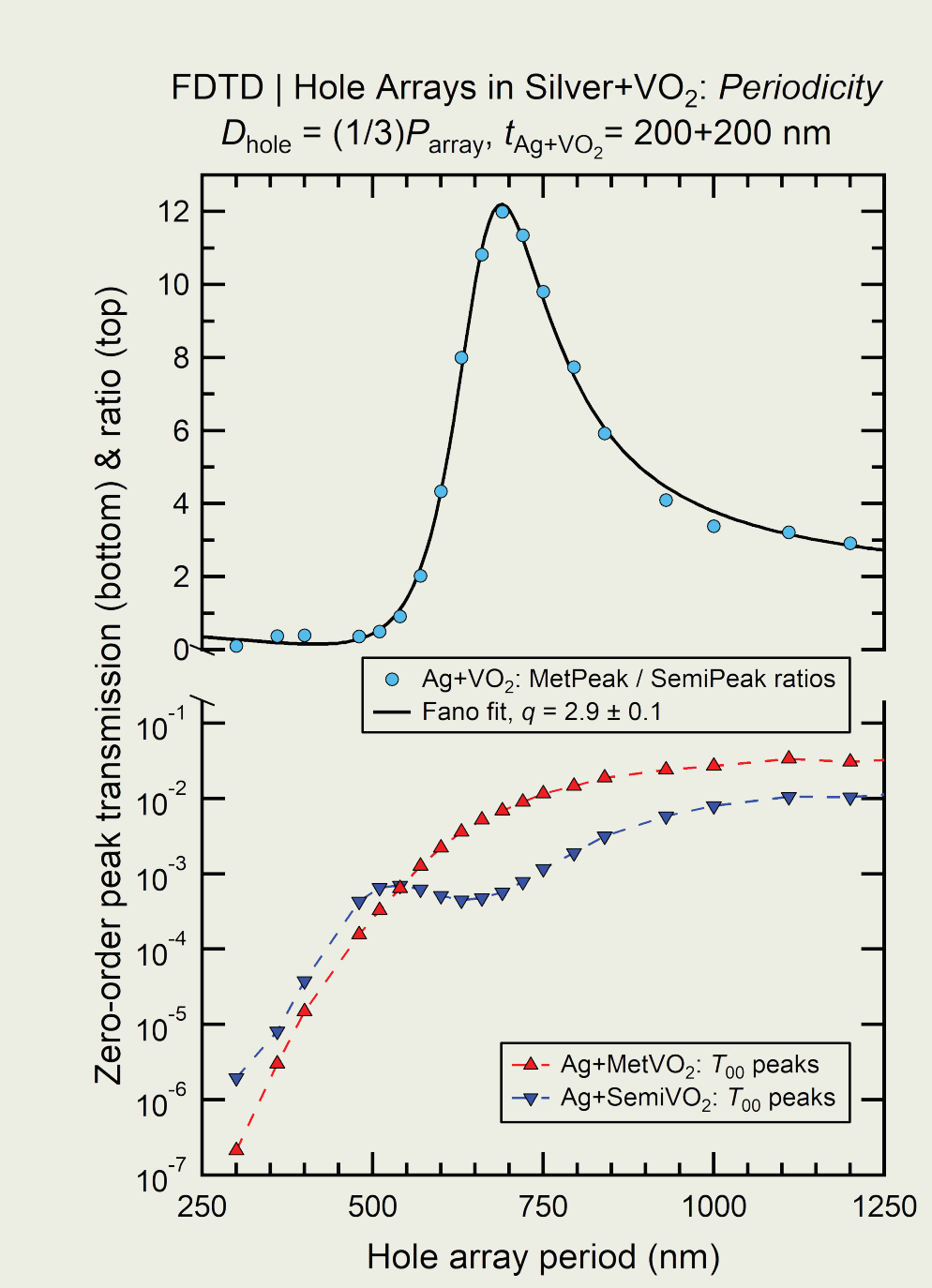
$$\text{FanoFit}(\lambda) = A + B \left[1 + \frac{q \Gamma_{\text{peak}}}{2(\lambda - \lambda_{\text{peak}})} \right]^2 \left/ \left[1 + \left(\frac{\Gamma_{\text{peak}}}{2(\lambda - \lambda_{\text{peak}})} \right)^2 \right] \right.$$

where *q* is the Fano asymmetry parameter related to the ratio of resonant and non-resonant contributions to the spectral profile.



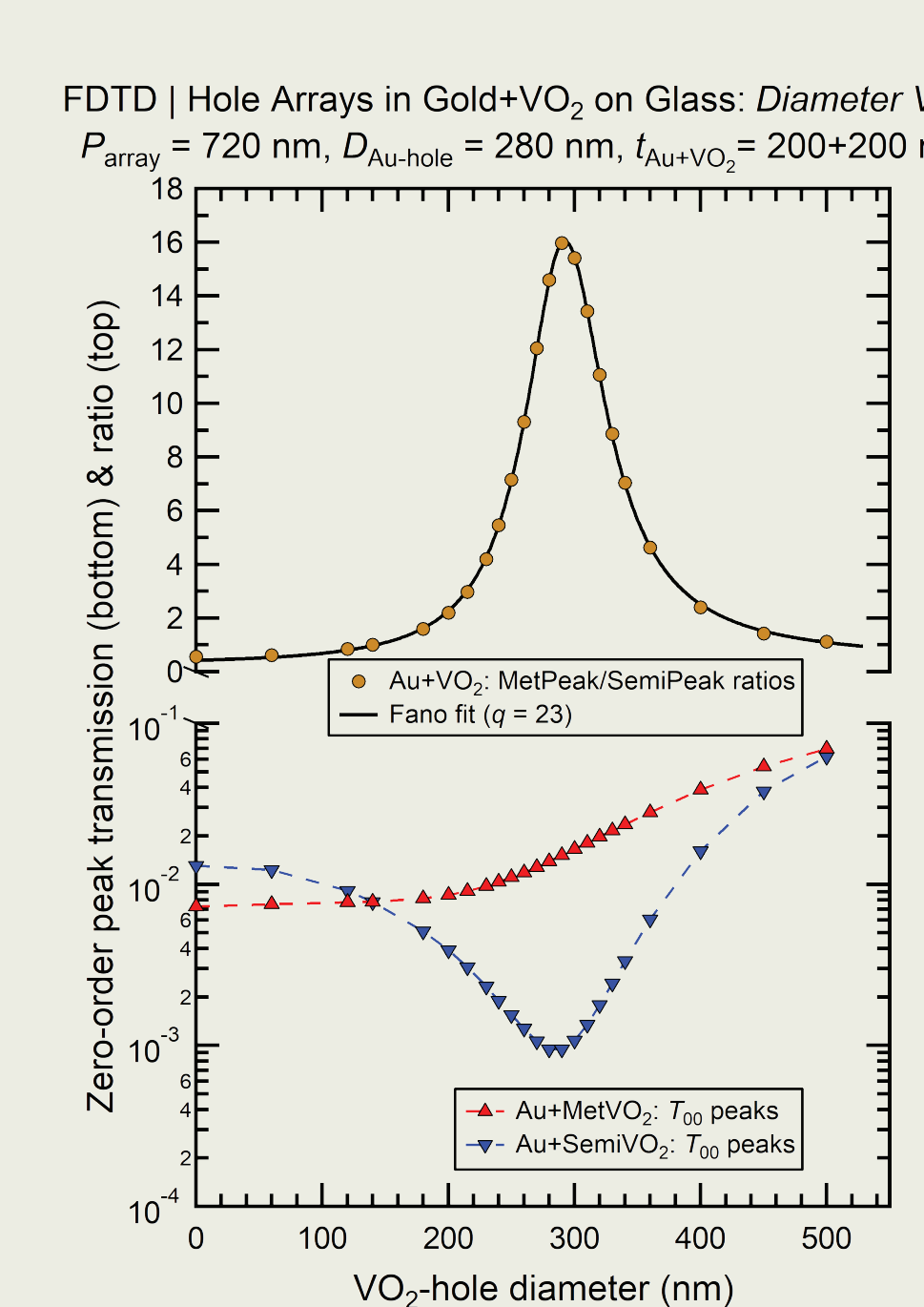
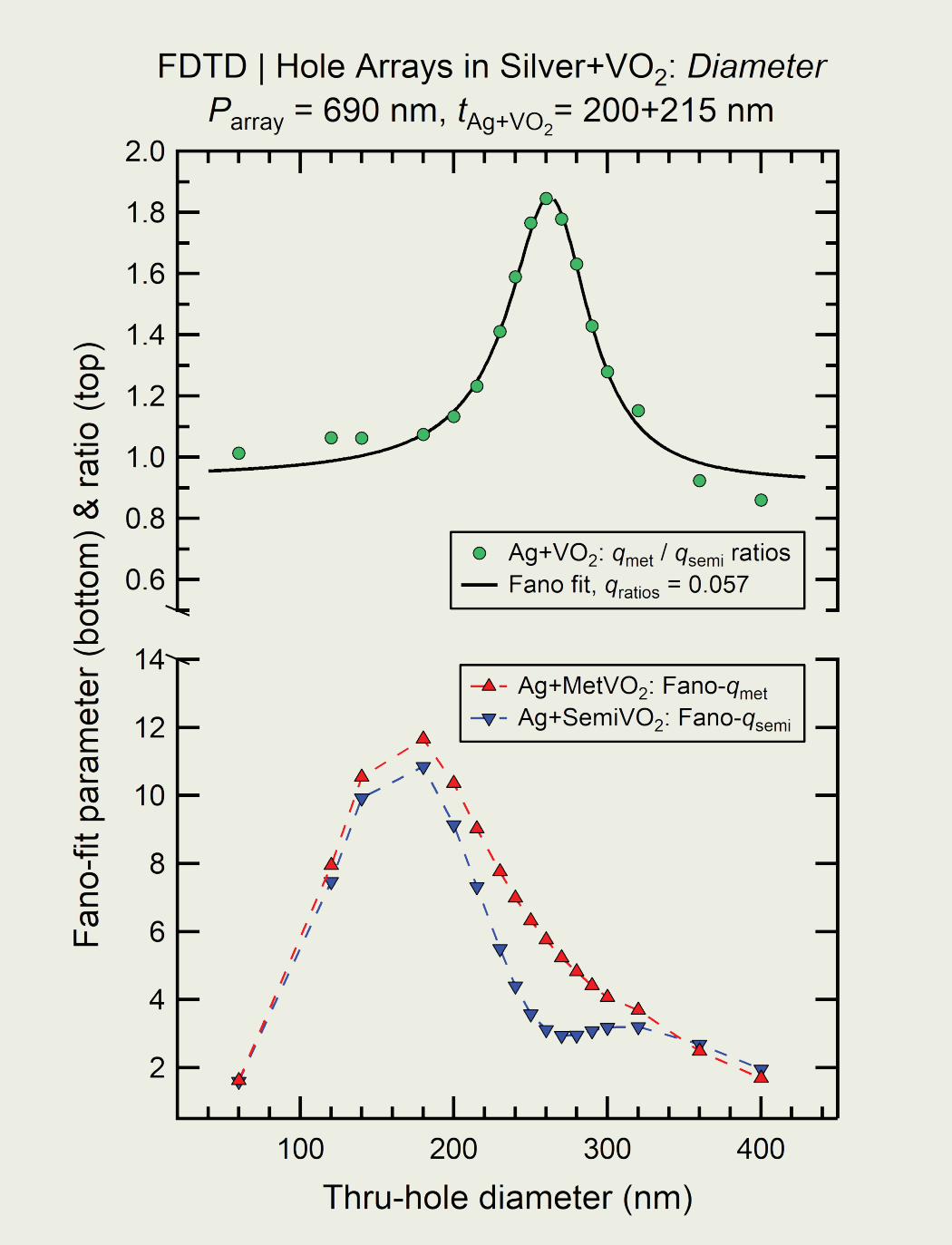
VARYING ARRAY PERIODICITY

As the period of the hole array in Au+VO₂ or Ag+VO₂ increases, the ratio of the EOT peak transmission *T*₀₀ in the metallic phase (MetVO₂) rises sharply and then decreases relative to the semiconducting phase (SemiVO₂). The images of the electric-field intensity and Poynting vector in YZ plane (orthogonal to the incident polarization) show the extent of penetration of the evanescent waves into the VO₂ layer. The peak ratios as a function of hole period follows very closely a Fano-type profile with either plasmonic metal, with or without a substrate. Intriguingly (see left figure for Ag+PlainVO₂), when the holes perforate only the plasmonic metal but not the VO₂ layer, the peak transmission in the semiconducting phase exceed that in the metallic phase for all array periods.



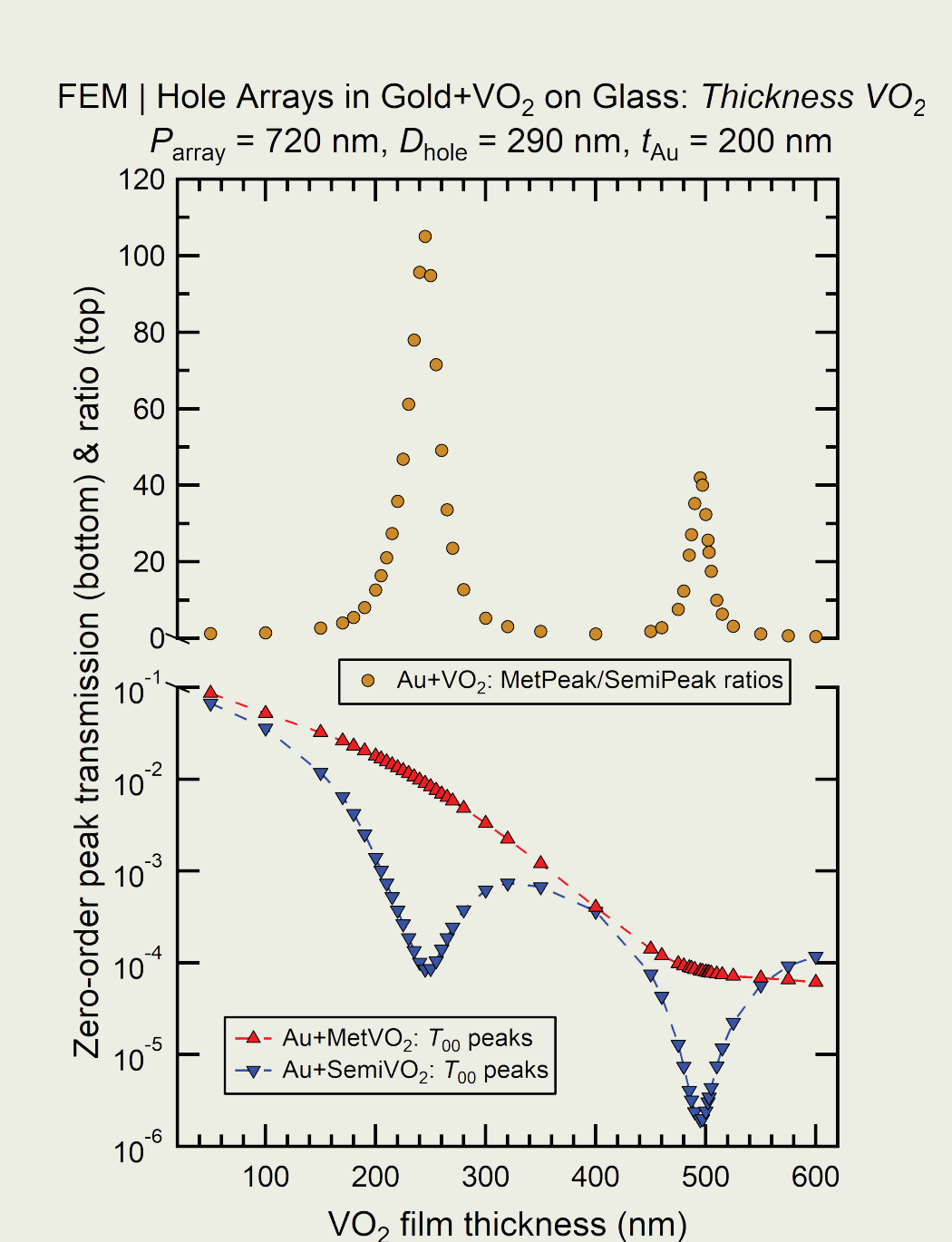
VARYING THRU-HOLE DIAMETER

Increasing the diameter of the holes that perforate both the plasmonic metal (Au or Ag) and VO₂ layers results in higher, broader and redshifted *T*₀₀ peaks, as expected. What is unexpected is that the ratio of the peak transmission in the metallic phase (MetVO₂) relative to that in the semiconducting phase (SemiVO₂) itself peaks for a hole diameter around 250 nm, before falling off toward unity. This diameter corresponds to a dip in the SemiVO₂ transmission. A Fano profile fits well to the *T*₀₀ ratio, though with a weak asymmetry. However, the ratio of Fano parameters *q* from the fits to the individual *T*₀₀ spectra has a more asymmetric shape; the corresponding “meta-Fano” fit (i.e., a Fano fit to Fano-*q* ratios) is also shown.

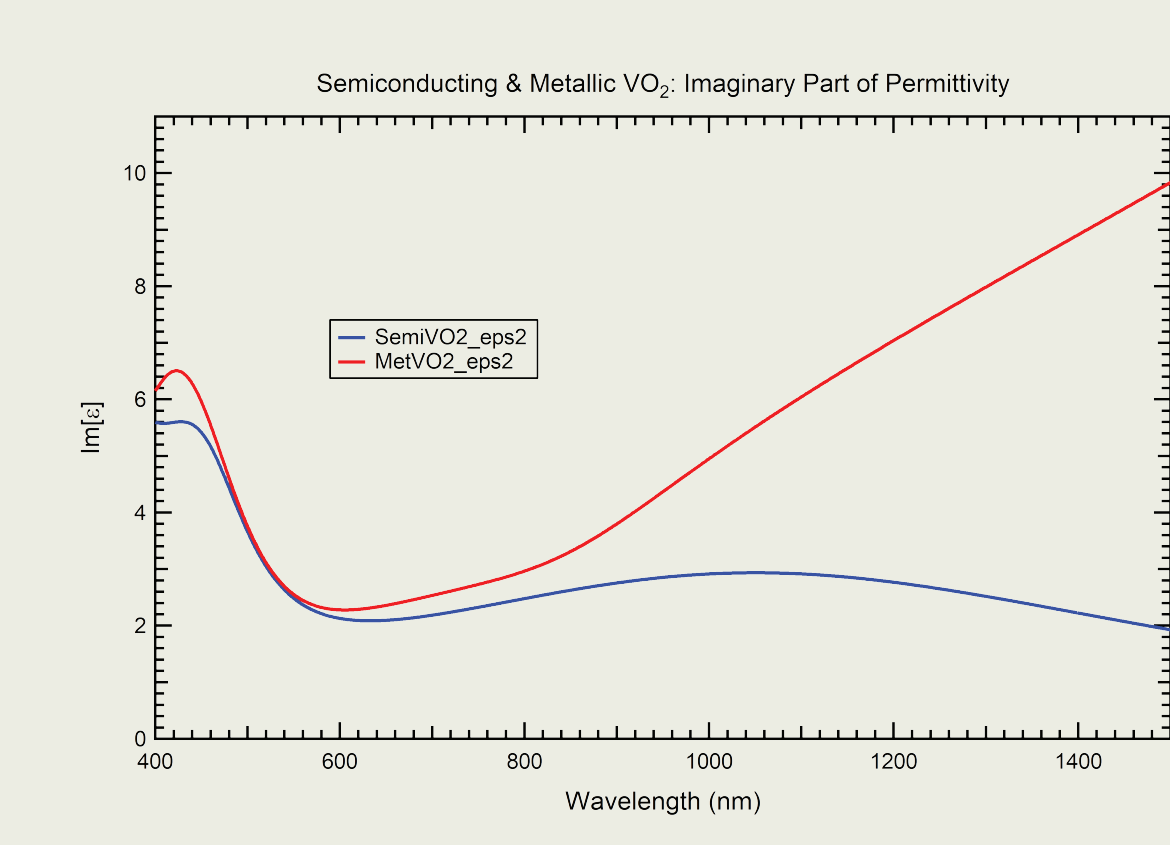
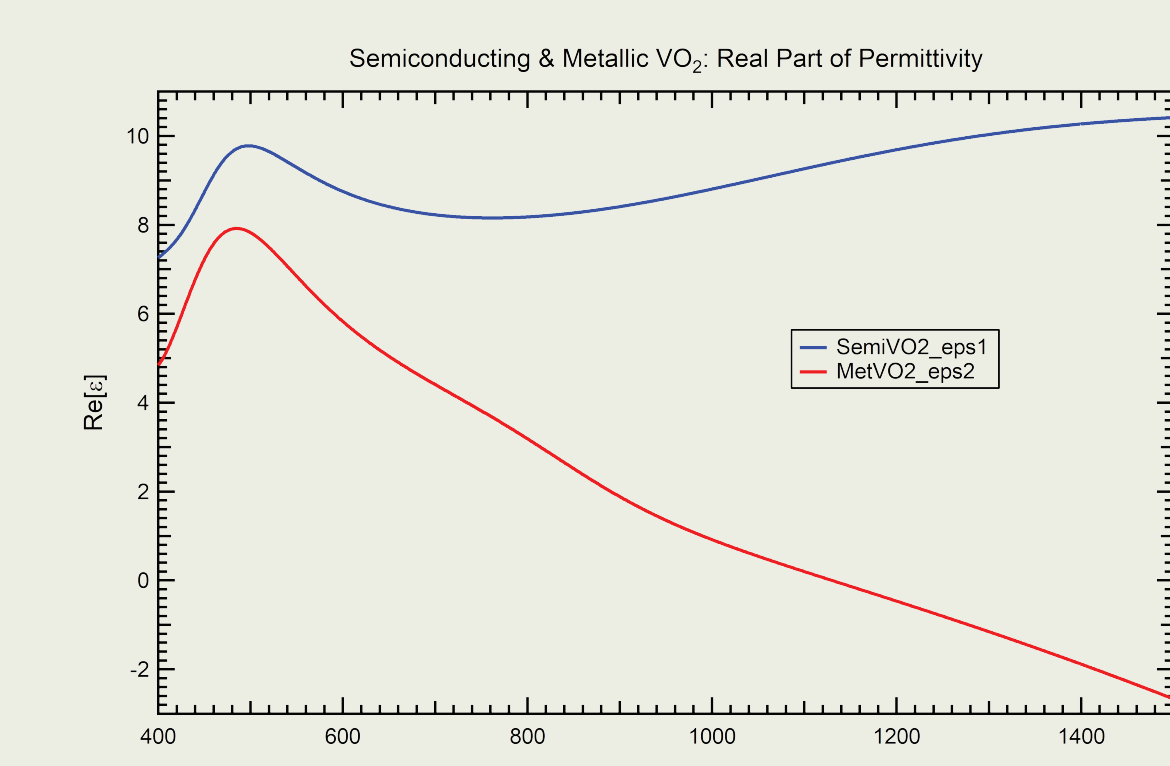


VARYING VO₂ HOLE DIAMETER OR FILM THICKNESS

A similar barely-asymmetric Fano profile fits the MetVO₂-to-SemiVO₂ ratio of *T*₀₀ peaks for a fixed hole diameter in the plasmonic layer and a variable hole diameter in the VO₂ layer. In this case the peaks broaden and redshift much less than in the thru-hole case because the spectral profiles are largely determined by the Au or Ag (fixed period and diameter) due to the EOT effect, whereas the VO₂ modulates mostly the transmission intensity. For the larger VO₂ holes, light emerging from the Au/Ag holes does not reach deep into the VO₂ material, so the effect of the different optical constants of the two VO₂ phases is relatively small, hence the ratio approaches unity.



When the thickness of the VO₂ layer is increased from zero, the *T*₀₀ ratio exhibits sharp symmetric peaks of diminishing amplitude, indicative of lossy Fabry-Perot resonances. Once again, as with variations in the array period and hole diameter, it is the semiconducting phase that behaves differently: The transmission dips in a narrow range around a specific thickness/period/diameter.



WHAT'S GOING ON?

Of course, the VO₂-induced modulation in the EOT through Au/Ag+VO₂ hole arrays can be traced back to the markedly different optical constants of the two VO₂ phases. The FDTD and FEM electromagnetic simulations reveal intriguing transmission trends—“reverse switching” of *T*₀₀ peaks and Fano-type profiles of peak ratios—but an explanation based on physical intuition remains elusive. We “wave hands” that the optimal parameters for maximizing the modulation are those for which the penetration of evanescent waves into the VO₂ film is least/most impeded by dissipation in the metallic/semiconducting phase. The goal of the CNMS portion of this project is to fabricate and test some of the simulated hole arrays to determine whether the Fano-type trends that seem so sensitive to the geometric parameters can survive under real-world experimental conditions (e.g., conical hole shape, surface roughness, and different optical constants). This work is currently underway at CNMS.

ACKNOWLEDGEMENTS

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