# METAL–SEMICONDUCTOR TRANSITIONS IN NANOSCALE VANADIUM DIOXIDE—THIN FILMS, SUBWAVELENGTH HOLES, AND NANOPARTICLES

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### APPENDIX A

### VANADIUM SESQUIOXIDE (V<sub>2</sub>O<sub>3</sub>) THIN FILMS

### Abstract

The following information is meant to provide some measure of continuity in our effort to make another material that undergoes a spectacular metal-insulator phase transition: vanadium sesquioxide,  $V_2O_3$ . A simple fabrication protocol is presented, along with some of the characterization data used to establish it. The two main steps are: (i) room-temperature PLD of amorphous  $V_xO_y$  with excess oxygen content, followed by (ii) high-temperature annealing in a reducing atmosphere and crystallization to stoichiometric  $V_2O_3$ .

## 1.1 Introduction

V<sub>2</sub>O<sub>3</sub> is the other famous oxide of vanadium. First discovered by Föex in 1946, vanadium sesquioxide undergoes a metal-to-insulator transition upon cooling and the reverse transition upon heating through  $T_c \approx 150$  K, accompanied by magnetic and crystallographic changes: at room temperature V<sub>2</sub>O<sub>3</sub> is a paramagnetic metal with a rhombohedral (corundum) lattice, while the low-temperature phase is an antiferromagnetic monoclinic insulator with a 0.6-eV bandgap.<sup>210</sup> Resistivity jumps by up to 6–7 orders of magnitude,<sup>21,211</sup> infrared (IR) transmission also increases,<sup>212–215</sup> and the first-order nature of the transition gives rise to hysteresis. The V<sub>2</sub>O<sub>3</sub> phase transition has long been regarded as a model for the Mott-Hubbard transition mechanism<sup>59,210,216,217</sup> (see also Section 1.1.3), with the monoclinic lattice distortion thought to originate simply from magnetostrictive forces<sup>63</sup> due to the peculiar magnetic ordering<sup>23</sup> below  $T_c$ . Unlike the structural transition of VO<sub>2</sub>, the change in crystal structure of V<sub>2</sub>O<sub>3</sub> is *not* associated with a Peierls instability and unit-cell doubling.<sup>59</sup> However, recent X-ray absorption measurements<sup>218</sup> of the temperature dependence of the local structure have identified a structural *precursor* to the metal-to-insulator transition of  $V_2O_3$ , namely a continuous increase in the monoclinic tilt starting well before the onset of the electronic and magnetic transition. These findings suggest that it may be the orbital degrees of freedom that drive the metal-to-insulator transition via changes in hybridization, which are in turn triggered by the monoclinic distortion.<sup>218</sup> Here too, as in the case of VO<sub>2</sub>, the web of cause-and-effect links in the transition mechanism of  $V_2O_3$  is yet to be untangled.

Studies devoted to bulk single crystals or thin films of  $V_2O_3$  abound in the scientific literature: according to Imada *et al.*,<sup>23</sup> more than 500 papers on  $V_2O_3$  had been published as of 1998. Yet only a handful of articles<sup>219–225</sup> deal with nanocrystalline  $V_2O_3$ —all reporting exclusively on chemical synthesis and/or catalysis; I know of no studies on the phase transition of  $V_2O_3$  nanoparticles (NPs). Observation of the optical switching of  $V_2O_3$  NPs would be more than a scientific "first". Studying nanostructures of another system that behaves like  $VO_2$  in terms of optical changes across the phase transition, and yet differs from  $VO_2$  as regards the roles of the various degrees of freedom involved, can only improve our understanding of how one solid-state phase transforms into another at the nanoscale. In particular, experimental data from "switchable"  $V_2O_3$  NPs would provide a new set of benchmarks for testing the "potent defect" model (see Section 1.3) of the microscopic origins of such solid-solid transformations.

# 1.2 A $V_2O_3$ "recipe"

The following deposition and annealing steps were found to yield good-quality  $V_2O_3$  films (~100 nm) on fused-silica substrates:

• Clean substrate: solvents (TCE, acetone, methanol/IPA, deionized water) and/or UVozone treatment.

• Deposit amorphous  $V_x O_y$  ( $y/x \ge 1.5$ ) film by room-temperature PLD: V-metal or  $V_2 O_3$ pressed-powder target; 300-mJ pulses focused to ~ 0.1 cm<sup>2</sup> on target surface; 15-cm targetto-substrate distance; 3-5 mtorr  $O_2$ .

• Reduce and crystallize deposited film by annealing in tube furnace, one sample at a time: (i) introduce 1 atm of flowing {4% H<sub>2</sub> + 96% Ar} into vacuum-tight quartz tube, pre-evacuated to 1–5 mtorr (use oil filter with the roughing pump); (ii) ramp up furnace temperature to 600 °C at 30 °C/min; (iii) dwell at 600 °C for 60 min; (iv) turn off heaters and let system cool down to (near) room temperature under flowing gas mix.

When viewed in transmission against white light, a "switchable" V<sub>2</sub>O<sub>3</sub> film should look some shade of grey, depending on thickness, but without a tinge of green, brown, or yellow. The acid test, as it were, for good-quality  $V_2O_3$  material is the optical switching in the vicinity of 150 K: sharpness, contrast, and hysteresis of the IR transmission across the metal-insulator transition. The transmission setup used here comprises: (i) fiber-coupled light source, either an IR laser at  $\lambda = 1330$  nm or a white-light tungsten-halogen lamp; (ii) mechanical chopper connected to a lock-in amplifier through a frequency generator; (iii) beamsplitter and CDD camera for visual inspection of the interrogated area via reflected and scattered light from the sample surface; (iv) one focusing and one collection lowmagnification micro-objectives, with the sample in-between; (v) pinhole aperture ( $\sim 0.5$ mm) for stray-light rejection; (vi) InGaAs IR detector with responsivity (amp/watt) greater than 10 % for  $\lambda = 800-1700$  nm; (vii) lock-in amplifier, with input signals from the chopper and the detector, and output of amplified and filtered DC signal proportional to the intensity of the transmitted light. The sample is mounted on a thin copper plate in contact with the heating/cooling pad of a micro-refrigerator assembly, which cools on the principle of expansion of highly pressurized dry  $N_2$  gas inside a series of micrometer-sized capillaries (Joule-Thomson effect) and heats by means of a resistor coil. The assembly is housed in a small optical chamber, continuously pumped to maintain roughing vacuum. The sample temperature is scanned and maintained  $(\pm 0.05 \text{ K})$  using a controller unit, which supplies power to the resistive heater based on the feedback from a temperature sensor under the heating plate.

#### **1.3** Different annealing temperatures

The first three figures show the effect of annealing temperature on the IR transmission (Figures A.1 and A.3) and X-ray diffraction (XRD) (Figure A.2) of two sets of initially amorphous  $V_x O_y$  films on fused-silica (SiO<sub>2</sub>) substrates. Upon visual inspection of these data, 600 °C was chosen for further experiments based on somewhat subjective criteria: (i) it is at the lower end of the temperature window that yielded films with "good switching", hence reducing the possibility that the (eventual) closely-spaced  $V_x O_y$  NPs would diffuse towards one another and coalesce during the anneal; (ii) unlike the 925-°C or 850-°C films, the film annealed at 600 °C has (nearly) completed its transition into the insulating state by the lowest measured temperature (85 K); (iii) switching contrast is larger for the 65-nm-thick 600-°C film than for its 700-°C counterpart; (iv) 600 °C has been popular processing temperature in previous V<sub>2</sub>O<sub>3</sub> studies.<sup>212-214</sup>

### 1.4 Different annealing times and ramp rates

The interesting curve in Figure A.4 is the one for slow ramp-up (3 °C/min) to the final annealing temperature of 600 °C. That particular film likely became sub-stoichiometric  $(V_2O_{y<3})$  before it had had a chance to crystallize into  $V_2O_3$ . A ramp-up rate of 30 °C/min and a dwell time of 1 hr were chosen for subsequent anneals, since dwell times of 0–4 hr had all produced switchable  $V_2O_3$  films (Figures A.1 and A.4).

#### **1.5** Different PLD target materials

Heeding an apropos observation by Schuler *et al.*<sup>213</sup> that  $V_2O_3$  films grown by reactive e-beam evaporation from a V-metal target or from a sintered  $V_2O_3$  powder exhibited "a drastic difference in quality", the former being the better, we performed depositions from four different PLD targets: V-metal disc from Cerac, inc. (http://www.cerac.com) and pressed-powder discs with nominal compositions of  $V_2O_3$ ,  $VO_2$ , and  $V_2O_5$  from Vin Karola Instruments (http://www.vinkarola.com). While the V and  $V_2O_3$  targets did produce films with steeper transmission hystereses than those from the other two targets, the differences in quality were hardly drastic (see Figure A.8).





Figure A.1: Relative IR transmission as a function of temperature (heating part of hysteresis cycle) for 140-nm-thick V<sub>2</sub>O<sub>3</sub> films on fused silica, H<sub>2</sub>-annealed at the indicated temperatures: 1100, 1000, 925, 850, 700 °C. Dwell times: 2 hr, except for the 925-°C film, which was annealed for 4 hr. Illumination sources: white-light lamp for the 925-°C film, and IR laser ( $\lambda = 1330$  nm) for the rest. The data for each film are normalized to the highest transmission in the low-temperature phase; the curves are offset vertically relative to one another for clarity.



Figure A.2: Room-temperature X-ray diffraction (XRD)  $\theta$ -2 $\theta$  scans (Cu- $K_{\alpha}$ ,  $\lambda = 1.54$  Å), for the above 140-nm-thick V<sub>2</sub>O<sub>3</sub> films on fused silica, H<sub>2</sub>-annealed at the indicated temperatures: (a) 700 °C; (b) 850 °C; (d) 925 °C; (e) 1000 °C; (f) 1100 °C. Part (c) shows the XRD scan for the film in (d) as deposited, *i.e.*, before annealing at 925 °C. Peaks at the powder-diffraction values (PDF #34-0187) of  $2\theta = 36.23^{\circ}$ , 65.193°, and 76.914° correspond to reflections from V<sub>2</sub>O<sub>3</sub> planes (1 1 0), (3 0 0), and (2 2 0), respectively.



 $V_2O_3$  on silica: Different anneal temperatures 2

Figure A.3: Relative IR transmission as a function of temperature for 65-nm-thick  $V_2O_3$  films, H<sub>2</sub>-annealed at the indicated temperatures: 550, 600, 650, 700, 750, 800 °C. Dwell time: 1 hr. Illumination: white-light lamp. Data for each film are normalized to the highest transmission in the low-temperature phase; the curves are offset vertically relative to one another for clarity. Filled symbols (red lines) correspond to the heating part of the hysteresis cycle, and open symbols (blue lines) to the cooling part.



Figure A.4: Relative IR transmission as a function of temperature for 65-nm-thick  $V_2O_3$  films (from the same initial batch as the films in Figure A.3), H<sub>2</sub>-annealed by ramping up the temperature to 600 °C at the indicated rates and holding it constant thereafter for the indicated dwell times. Illumination source: white-light lamp. Data for each film are normalized to the highest transmission in the low-temperature phase; the curves are offset vertically relative to one another for clarity. Filled symbols (red lines) correspond to the heating part of the hysteresis cycle, and open symbols (blue lines) to the cooling part.



Figure A.5: Example of experimental (circles) and simulated (lines) RBS spectra for vanadium-oxide film on fused-silica (SiO<sub>2</sub>) substrate: (a) before H<sub>2</sub> anneal (*i.e.*, as deposited); (b) after H<sub>2</sub> anneal for 1 hr at 600 °C. The film was grown by room-temperature PLD from a V<sub>2</sub>O<sub>3</sub> target in 3 mtorr O<sub>2</sub> background gas. Simulations were performed using the SIMNRA program.<sup>107</sup> Analysis: (a) V<sub>2</sub>O<sub>3.12±0.01</sub>,  $t \approx 101$  nm; (b) V<sub>2</sub>O<sub>3.03±0.01</sub>,  $t \approx 99$  nm.



Figure A.6: RBS-measured stoichiometry (see example in Figure A.5), before and after H<sub>2</sub> anneal (1 hr at 600 °C), for vanadium-oxide films deposited from V-metal (in 3 mtorr  $O_2$ ),  $V_2O_3$  (in 3 mtorr  $O_2$ ),  $VO_2$  (in vacuum), and  $V_2O_5$  (in vacuum) PLD targets. The number above each bar denotes the thickness in nanometers.



Figure A.7: Room-temperature XRD scans (Cu– $K_{\alpha}$ ,  $\lambda = 1.54$  Å), before and after H<sub>2</sub> anneal (1 hr at 600 °C), for vanadium-oxide films deposited from (a) V-metal (in 3 mtorr O<sub>2</sub>), (b) V<sub>2</sub>O<sub>3</sub> (in 3 mtorr O<sub>2</sub>), (c) VO<sub>2</sub> (in vacuum), and (d) V<sub>2</sub>O<sub>5</sub> (in vacuum) PLD targets. Peaks at the powder-diffraction values (PDF #34-0187) of  $2\theta = 36.23^{\circ}$  and  $65.193^{\circ}$  correspond to reflections from V<sub>2</sub>O<sub>3</sub> planes (1 1 0) and (3 0 0), respectively.



Hystereses of V<sub>2</sub>O<sub>3</sub> on silica: Different PLD targets

**Figure A.8:** Optical transmittance for H<sub>2</sub>-annealed  $V_2O_3$  films deposited from different PLD target materials. Anneal: 1 hr at 600 °C. Illumination source: white-light lamp. Data for each film are normalized at each temperature point to transmission through the bare fused-silica substrate; the curves are offset vertically relative to one another for clarity. Filled symbols (red lines) correspond to the heating part of the hysteresis cycle, and open symbols (blue lines) to the cooling part.

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